QUANTUM MECHANICS Appendix 3

Enrico Iacopini

QUANTUM MECHANICS Appendix 3 Finite square well: Wave function normalization

Enrico Iacopini

October 16, 2019

Enrico Iacopini

QUANTUM MECHANICS Appendix 3

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >

October 16, 2019 1 / 9

3

Let us start from the even solutions that we have obtained for the (bound) stationary states associated to the finite square well. We have found

$$\begin{aligned} |x| \ge a &: \ \psi(x) = \alpha \, e^{-r|x|} \\ x < a &: \ \psi(x) = B \cos(kx) \end{aligned}$$

where $\alpha = B e^{\eta} \cos \xi$ and $\eta = a r = \xi t g \xi$.

2 We intend to determine the constant *B* for which $|\psi|^2$ is normalized (*B* is unique, a part a possible complex phase ...).

イロト 不得 トイヨト イヨト ニヨー

Let us start from the even solutions that we have obtained for the (bound) stationary states associated to the finite square well. We have found

$$\begin{aligned} |x| &\ge a : \ \psi(x) = \alpha \, e^{-r|x|} \\ x &< a : \ \psi(x) = B \cos(kx) \end{aligned}$$

where $\alpha = B e^{\eta} \cos \xi$ and $\eta = a r = \xi t g \xi$.

We intend to determine the constant B for which $|\psi|^2$ is normalized (B is unique, a part a possible complex phase ...).

Let us start from the even solutions that we have obtained for the (bound) stationary states associated to the finite square well. We have found

$$\begin{aligned} |x| &\ge a : \ \psi(x) = \alpha \, e^{-r|x|} \\ x &< a : \ \psi(x) = B \cos(kx) \end{aligned}$$

where $\alpha = B e^{\eta} \cos \xi$ and $\eta = a r = \xi t g \xi$.

We intend to determine the constant *B* for which $|\psi|^2$ is normalized (*B* is unique, a part a possible complex phase ...).

Since $|\psi|^2$ is even, we will integrate only for x > 0 and then double the result. We have

$$\int_{0}^{a} dx |\psi(x)|^{2} = B^{2} \int_{0}^{a} dx \cos^{2}(kx) =$$

$$= \frac{1}{2}B^{2} \int_{0}^{a} dx [1 + \cos(2kx)] =$$

$$= B^{2} \left[\frac{a}{2} + \frac{1}{2}\frac{\sin(2ka)}{2k}\right] =$$

$$= B^{2} \frac{2ak + \sin(2ak)}{4k} = B^{2} \frac{a}{2} \frac{2\xi + \sin(2\xi)}{2\xi} =$$

$$= B^{2} \frac{a}{2} \left(1 + \frac{\sin\xi\cos\xi tg\xi}{\xi tg\xi}\right) =$$

$$= B^{2} \frac{a}{2} \left(1 + \frac{\sin^{2}\xi}{\xi tg\xi}\right)$$
Entro lacopia QUANTUM MECHANICS Appendix 3 October 16.

16, 2019 3 / 9

I From the region in which x > a, we get

$$\int_{a}^{\infty} dx |\psi(x)|^{2} = \alpha^{2} \int_{a}^{\infty} dx e^{-2rx} =$$
$$= \alpha^{2} \left(-\frac{1}{2r}\right) e^{-2rx} \Big|_{a}^{\infty} =$$
$$= \alpha^{2} \frac{1}{2r} e^{-2ar} = a \alpha^{2} \frac{e^{-2\eta}}{2\eta}$$

2 But

$$\alpha = B e^{\eta} \cos\xi \Rightarrow \int_{a}^{\infty} dx |\psi(x)|^{2} =$$
$$= a B^{2} e^{2\eta} \cos^{2}\xi \frac{e^{-2\eta}}{2\eta} = \frac{a}{2} B^{2} \frac{\cos^{2}\xi}{\eta}$$

QUANTUM MECHANICS Appendix 3

3

• From the region in which x > a, we get

$$\int_{a}^{\infty} dx |\psi(x)|^{2} = \alpha^{2} \int_{a}^{\infty} dx e^{-2rx} =$$
$$= \alpha^{2} \left(-\frac{1}{2r}\right) e^{-2rx} \Big|_{a}^{\infty} =$$
$$= \alpha^{2} \frac{1}{2r} e^{-2ar} = a \alpha^{2} \frac{e^{-2\eta}}{2\eta}$$

2 But

$$\alpha = B e^{\eta} \cos\xi \Rightarrow \int_{a}^{\infty} dx |\psi(x)|^{2} =$$
$$= a B^{2} e^{2\eta} \cos^{2}\xi \frac{e^{-2\eta}}{2\eta} = \frac{a}{2} B^{2} \frac{\cos^{2}\xi}{\eta}$$

QUANTUM MECHANICS Appendix 3

QUANTUM MECHANICS Appendix 3

In conclusion, for the even solutions we have

$$\int_{-\infty}^{+\infty} dx \, |\psi(x)|^2 = a \, B^2 \left[\left(1 + \frac{\sin^2 \xi}{\eta} \right) + \frac{\cos^2 \xi}{\eta} \right] = a \, B^2 \left(1 + \frac{1}{\eta} \right)$$

and therefore

$$\int_{-\infty}^{+\infty} dx \, |\psi(x)|^2 = 1 \quad \Rightarrow$$
$$\Rightarrow \quad B^{-1} = \sqrt{a\left(1 + \frac{1}{\eta}\right)}$$

QUANTUM MECHANICS Appendix 3

Enrico Iacopini

3

Let us come, now, to consider the odd solutions, for which

$$\begin{array}{rcl} x \geq a & : & \psi(x) = \alpha \, e^{-rx} \\ x \leq -a & : & \psi(x) = -\alpha \, e^{rx} \\ x < a & : & \psi(x) = A \sin(kx) \end{array}$$

where $\alpha = A e^{\eta} \sin \xi$ and $\eta = kr = -\xi ctg\xi$.

QUANTUM MECHANICS Appendix 3

Enrico Iacopini

Since also in this case $|\psi|^2$ is **even**, we will integrate only for x > 0 and then double the result. We have

$$\int_{0}^{a} dx |\psi(x)|^{2} = A^{2} \int_{0}^{a} dx \sin^{2}(kx) =$$

$$= \frac{1}{2} A^{2} \int_{0}^{a} dx [1 - \cos(2kx)] =$$

$$= A^{2} \left[\frac{a}{2} - \frac{1}{2} \frac{\sin(2ka)}{2k} \right] = A^{2} \frac{a}{2} \left(1 - \frac{2\sin\xi\cos\xi}{2\xi} \right) =$$

$$= A^{2} \frac{a}{2} \left(1 - \frac{\sin\xi\cos\xi(-\operatorname{ctg}\xi)}{\xi(-\operatorname{ctg}\xi)} \right) =$$

$$= A^{2} \frac{a}{2} \left(1 + \frac{\cos^{2}\xi}{\eta} \right)$$

QUANTUM MECHANICS Appendix 3

1 From the x > a region, we get again

$$\int_{a}^{\infty} dx |\psi(x)|^{2} = \alpha^{2} \int_{a}^{\infty} dx e^{-2rx} =$$
$$= \alpha^{2} \left(-\frac{1}{2r}\right) e^{-2rx} \Big|_{a}^{\infty} =$$
$$= \alpha^{2} \frac{1}{2r} e^{-2ar} = a \alpha^{2} \frac{e^{-2\eta}}{2\eta}$$

) But nov

$$\alpha = A e^{\eta} \sin\xi \Rightarrow \int_{a}^{\infty} dx |\psi(x)|^{2} =$$
$$= a A^{2} e^{2\eta} \sin^{2}\xi \frac{e^{-2\eta}}{2\eta} = \frac{a}{2} A^{2} \frac{\sin^{2}\xi}{\eta}$$

Enrico Iacopini

QUANTUM MECHANICS Appendix 3

イロト 不得 トイヨト イヨト ニヨー

If the x > a region, we get again

$$\int_{a}^{\infty} dx |\psi(x)|^{2} = \alpha^{2} \int_{a}^{\infty} dx e^{-2rx} =$$
$$= \alpha^{2} \left(-\frac{1}{2r}\right) e^{-2rx} \Big|_{a}^{\infty} =$$
$$= \alpha^{2} \frac{1}{2r} e^{-2ar} = a \alpha^{2} \frac{e^{-2\eta}}{2\eta}$$

2 But now

$$\alpha = A e^{\eta} \sin \xi \Rightarrow \int_{a}^{\infty} dx |\psi(x)|^{2} =$$
$$= a A^{2} e^{2\eta} \sin^{2} \xi \frac{e^{-2\eta}}{2\eta} = \frac{a}{2} A^{2} \frac{\sin^{2} \xi}{\eta}$$

In conclusion, also for the odd solutions we have

$$\int_{-\infty}^{+\infty} dx \, |\psi(x)|^2 = a \, A^2 \, \left[\left(1 + \frac{\cos^2 \xi}{\eta} \right) + \frac{\sin^2 \xi}{\eta} \right] = \\ = a \, A^2 \, \left(1 + \frac{1}{\eta} \right)$$

and therefore, again

$$\int_{-\infty}^{+\infty} dx \, |\psi(x)|^2 = 1 \quad \Rightarrow$$
$$\Rightarrow \quad A^{-1} = \sqrt{a\left(1 + \frac{1}{\eta}\right)}$$

Enrico Iacopini

QUANTUM MECHANICS Appendix 3

QUANTUM MECHANICS Appendix 3